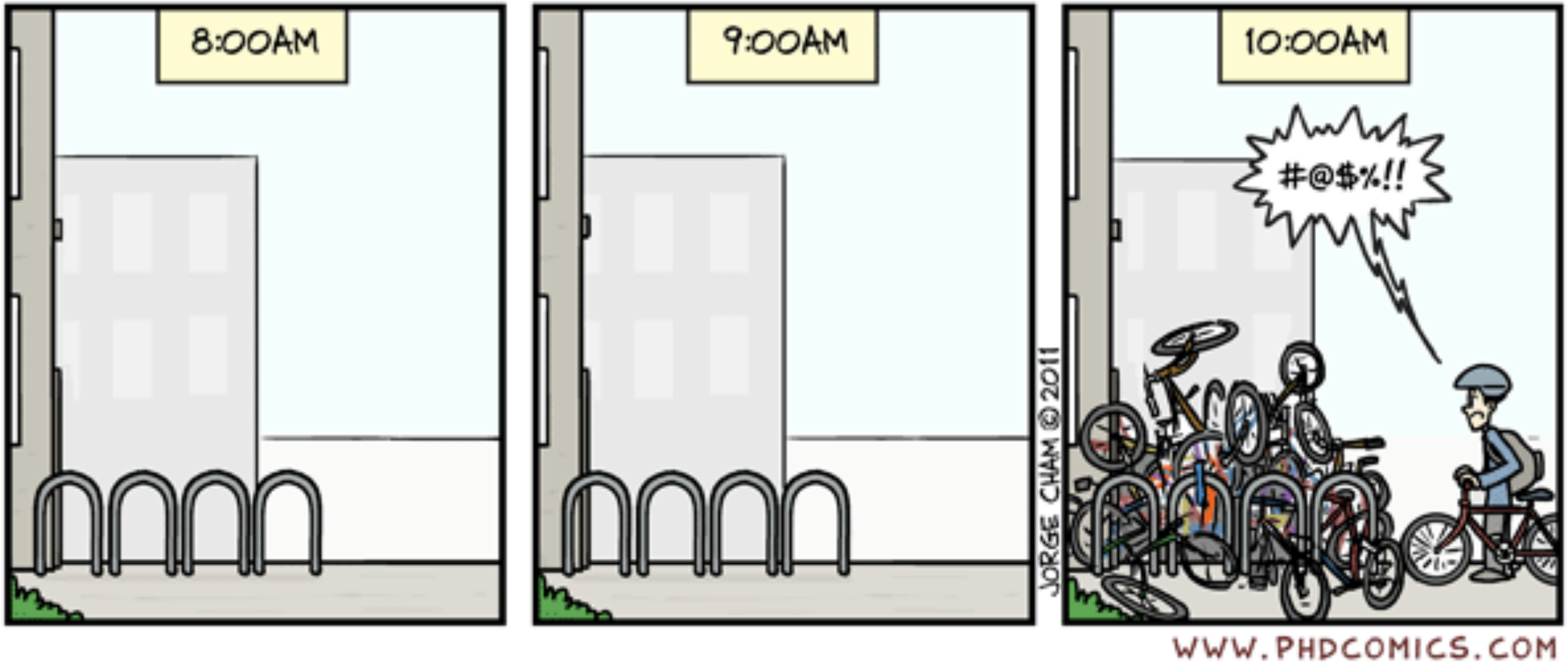
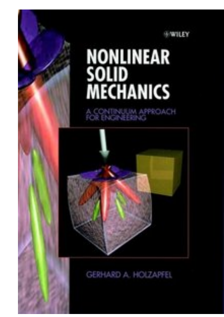
07 - balance principles



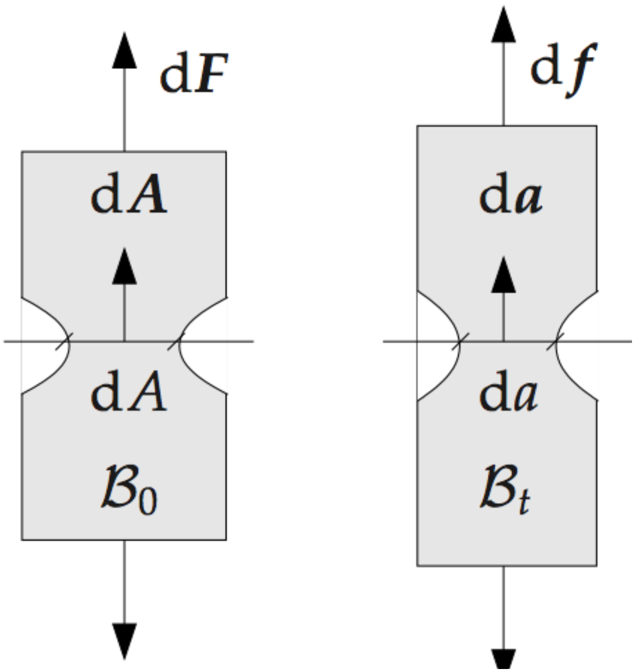
holzapfel 'nonlinear solid mechanics' [2000], chapter 4, pages 131-161

me338 - syllabus

day	date		topic	chapters	pages
tue	sep	25	why continuum mechanics?		
thu	sep	27	introduction to vectors and tensors	1.1-1.5	1-32
tue	oct	02	introduction to vectors and tensors	1.6-1.9	32-55
thu	oct	04	kinematics	2.1-2.4	55-76
tue	oct	09	kinematics	2.5-2.8	76-109
thu	oct	11	concept of stress	3.1-3.4	109-131
tue	oct	16	balance principles	4.1-4.4	131-161
thu	oct	18	balance principles	4.5-4.7	161-179
tue	oct	23	aspects of objectivity	5.1-5.4	179-205
thu	oct	25	hyperelastic materials	6.1-6.2	205-222
tue	oct	30	hyperelastic materials	6.3-6.5	222-252
thu	nov	01	hyperelastic materials	6.6-6.8	252-278
tue	nov	06	hyperelastic materials	6.9-6.11	278-305
thu	nov	08	thermodynamics of materials	7.1-7.6	305-337
tue	nov	13	midterm prep		
thu	nov	15	midterm		
tue	nov	27	thermodynamics of materials	7.7-7.9	337-371
thu	nov	29	variational principles	8.1-8.3	371-392
tue	dec	04	variational principles	8.4-8.6	392-414
thu	dec	06	final project discussion		



07 - balance principles



cauchy / true stress
relates spatial force to spatial area

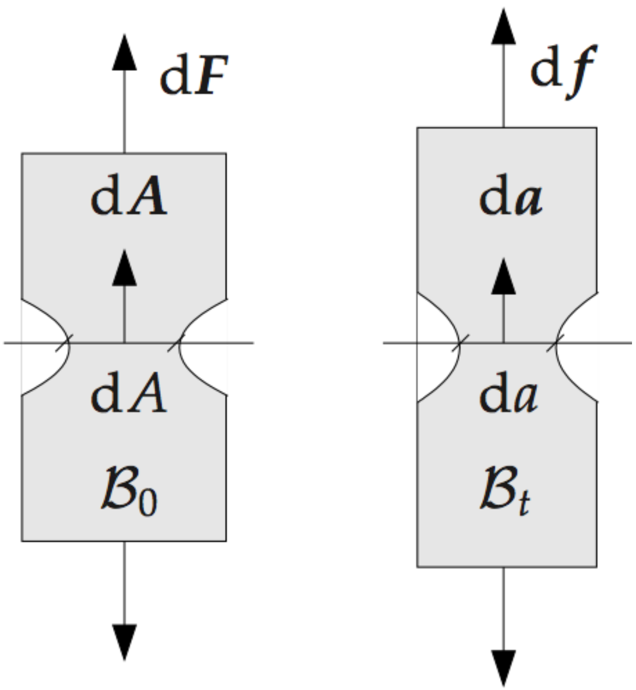
$$\mathrm{d}\boldsymbol{f} = \boldsymbol{t}\,\mathrm{d}a = \boldsymbol{\sigma}\cdot\boldsymbol{n}\,\mathrm{d}a = \boldsymbol{\sigma}\cdot\mathrm{d}\boldsymbol{a}$$

first piola kirchhoff / nominal stress
relates spatial force to material area

$$\mathrm{d}\boldsymbol{f} = \boldsymbol{t}\,\mathrm{d}a = \boldsymbol{\sigma}\cdot\boldsymbol{n}\,\mathrm{d}a = \boldsymbol{\sigma}\cdot\mathrm{d}\boldsymbol{a} = J\,\boldsymbol{\sigma}\cdot\boldsymbol{F}^{-\mathrm{t}}\cdot\mathrm{d}\boldsymbol{A} = \boldsymbol{P}\cdot\mathrm{d}\boldsymbol{A}$$

second piola kirchhoff stress
relates material force to material area

$$\mathrm{d}\boldsymbol{F} = \boldsymbol{F}^{-1}\cdot\mathrm{d}\boldsymbol{f} = \boldsymbol{F}^{-1}\cdot\boldsymbol{P}\cdot\mathrm{d}\boldsymbol{A} = J\,\boldsymbol{F}^{-1}\cdot\boldsymbol{\sigma}\cdot\boldsymbol{F}^{-\mathrm{t}}\cdot\mathrm{d}\boldsymbol{A} = \boldsymbol{S}\cdot\mathrm{d}\boldsymbol{A}$$



cauchy / true stress

relates spatial force to spatial area

$$\boldsymbol{t} = \boldsymbol{\sigma} \cdot \boldsymbol{n} \qquad t_i = \sigma_{ij}\, n_j$$

first piola kirchhoff / nominal stress

relates spatial force to material area

$$\boldsymbol{P} = J\, \boldsymbol{\sigma} \cdot \boldsymbol{F}^{-t} \qquad P_{iJ} = J\, \sigma_{ik}\, F^{-t}_{kJ}$$

second piola kirchhoff stress

relates material force to material area

$$\boldsymbol{S} = \boldsymbol{F}^{-1} \cdot \boldsymbol{P} = J\, \boldsymbol{F}^{-1} \cdot \boldsymbol{\sigma} \cdot \boldsymbol{F}^{-t}$$

$$S_{IJ} = F^{-1}_{Ik}\, P_{kJ} = J\, F^{-1}_{Ik}\, \sigma_{kl}\, F^{-t}_{lJ}$$

gustav robert kirchhoff
[1824-1887]

first piola kirchhoff

$$P = F \cdot S$$
$$P = J\, \sigma \cdot F^{-t}$$

augustin louis caucy
[1789-1857]

second piola kirchhoff

$$S = F^{-1} \cdot P$$
$$S = J\, F^{-1} \cdot \sigma \cdot F^{-t}$$

cauchy

$$\sigma = \frac{1}{J}\, P \cdot F^{t}$$
$$\sigma = \frac{1}{J}\, F \cdot S \cdot F^{t}$$

covariant / strains

$E \;= F^{t} \cdot e \cdot F$ ← pull back

$E_{IJ} = F^{t}_{Ik}\, e_{kl}\, F_{lJ}$

push forward →

$e \;= F^{-t} \cdot E \cdot F^{-1}$

$e_{ij} = F^{-t}_{iK}\, E_{KL}\, F^{-1}_{Lj}$

contravariant / stresses

$S \;= J\, F^{-1} \cdot \sigma \cdot F^{-t}$ ← pull back

$S_{IJ} = J\, F^{-1}_{Ik}\, \sigma_{kl}\, F^{-t}_{lJ}$

push forward →

$\sigma \;= \frac{1}{J} F \cdot S \cdot F^{t}$

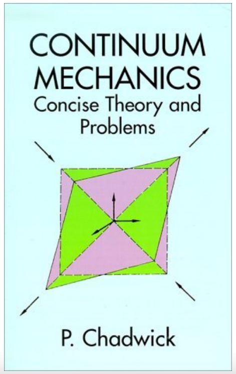
$\sigma_{ij} = \frac{1}{J} F_{iK}\, S_{KL}\, F^{t}_{Lj}$

balance equations

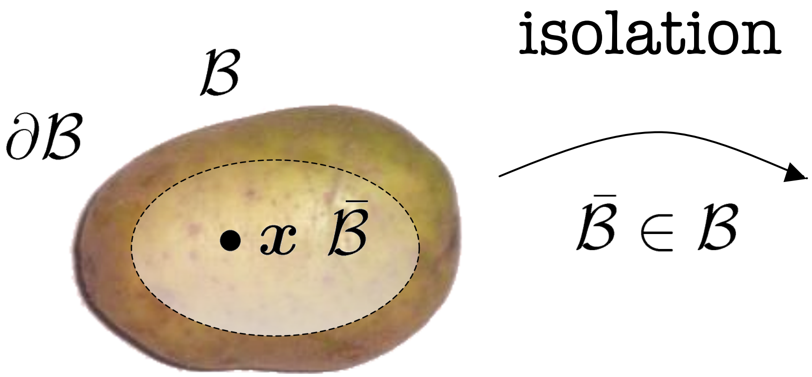
balance equations [ˈbæl.əns ɪˈkweɪ.ʒəns] of mass, momentum, angular momentum and energy, supplemented with an entropy inequality constitute the set of conservation laws. the law of **conservation of mass**/matter states that the **mass of a closed system** of substances will remain **constant**, regardless of the processes acting inside the system. the principle of conservation of momentum states that the total momentum of a closed system of objects is constant.

balance equations

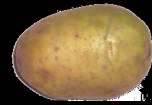
balance equations [ˈbæl.əns ɪˈkweɪ.ʒəns] of mass, linear momentum, angular momentum and energy **apply to all material bodies.** each one gives rise to a field equation, holing on the configurations of a body in a sufficiently smooth motion and a jump condition on surfaces of discontinuity. like position, time and body, the concepts of mass, force, heating and internal energy which enter into the formulation of the balance equations are regarded as having primitive status in continuum mechanics.

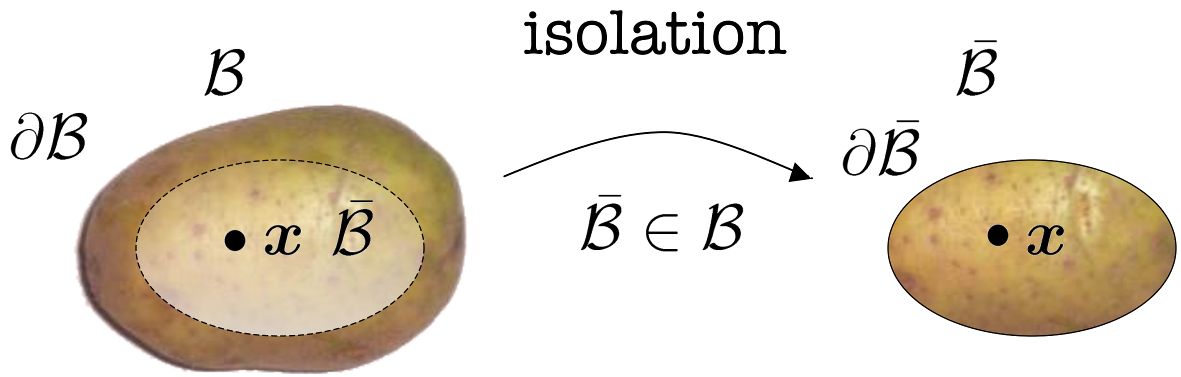


chadwick 'continuum mechanics' [1976]



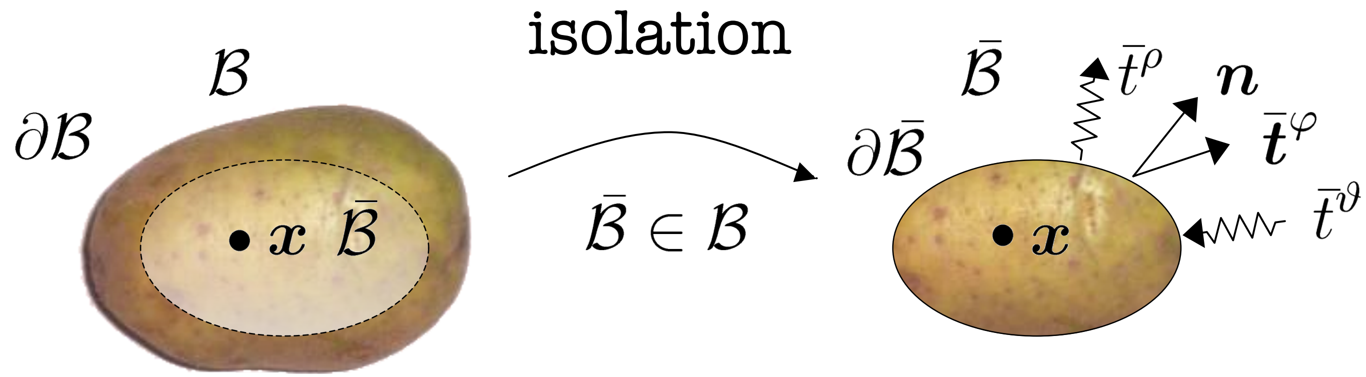
[1] isolate subset $\bar{\mathcal{B}}$ from $\mathcal{B}$

potato  balance equations



[1] isolate subset $\bar{\mathcal{B}}$ from $\mathcal{B}$

[2] **characterize** influence of remaining body through phenomenological quantities - contact fluxes $\bar{t}^{\rho}$, $\bar{\boldsymbol{t}}^{\varphi}$& $\bar{t}^{\vartheta}$

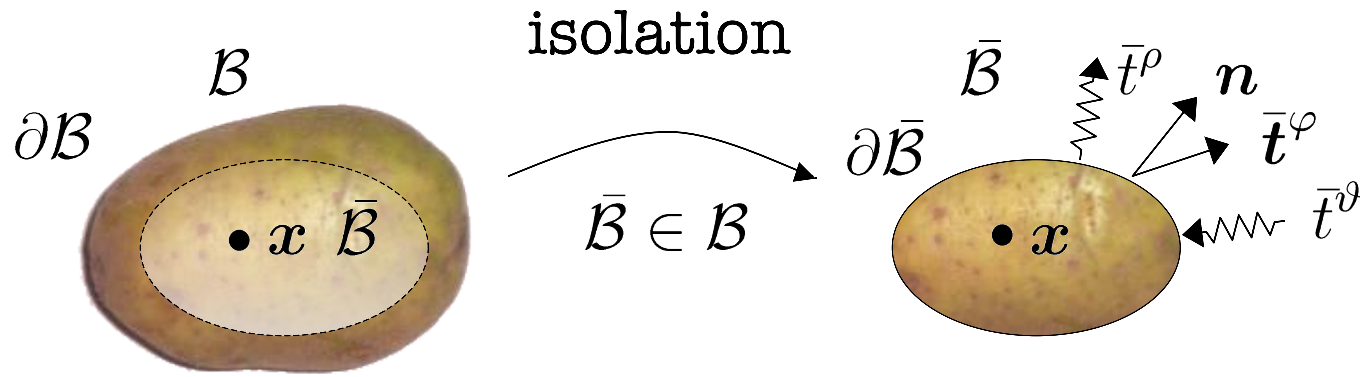


[1] isolate subset $\bar{\mathcal{B}}$ from $\mathcal{B}$

[2] characterize influence of remaining body through phenomenological quantities - contact fluxes $\bar{t}^{\rho}$, $\bar{\boldsymbol{t}}^{\varphi}$& $\bar{t}^{\vartheta}$

[3] **define** basic physical quantities - mass, linear and angular momentum, energy

potato  balance equations



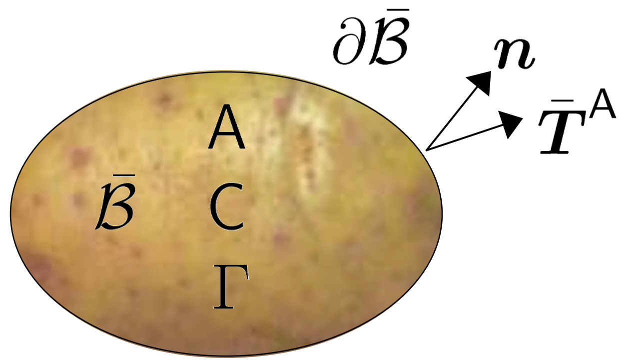
[1] isolate subset $\bar{\mathcal{B}}$ from $\mathcal{B}$

[2] characterize influence of remaining body through phenomenological quantities - contact fluxes $\bar{t}^{\rho}$, $\bar{\boldsymbol{t}}^{\varphi}$ & $\bar{t}^{\vartheta}$

[3] define basic physical quantities - mass, linear and angular momentum, energy

[4] **postulate** balance of these quantities

generic balance equation



general format

A ... balance quantity

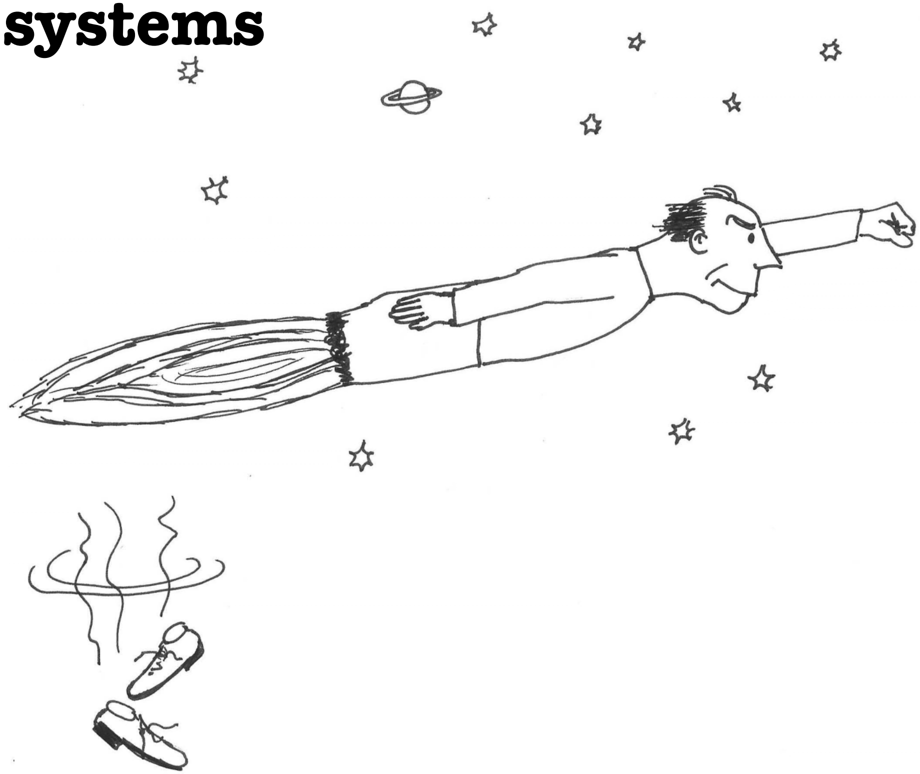
$\mathbf{B}$ ... flux $\quad \mathbf{B} \cdot \boldsymbol{n} = \bar{\boldsymbol{T}}^{\mathsf{A}}$

C ... source

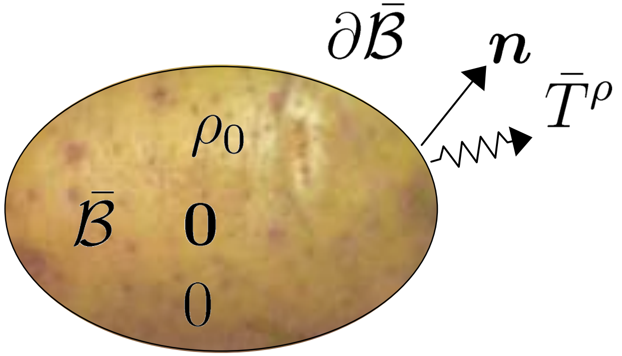
Γ ... production

$$\mathrm{D}_t \mathsf{A} = \mathrm{Div}(\mathbf{B}) + \mathsf{C} + \Gamma$$

- here: **closed systems**



- unlike open systems **closed systems** have a **constant mass**
- examples of open systems:
  rocket propulsion and **biological growth** (me337)



ρ_0 ... density

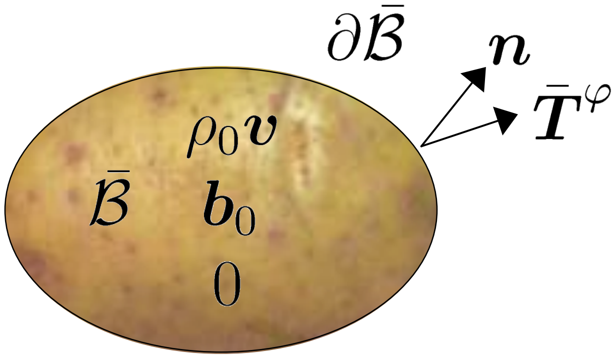
$\mathbf{0}$ ... no mass flux $\quad \bar{T}^\rho = 0$

0 ... no mass source

0 ... no mass production

continuity equation $\quad \mathrm{D}_t \rho_0 = 0$

balance of (linear) momentum



$\rho_0\,\boldsymbol{v}$ ... linear momentum density

$\boldsymbol{P}$ ... momentum flux - stress $\qquad \boldsymbol{P}\cdot\boldsymbol{n} = \bar{\boldsymbol{T}}^{\varphi}$

$\boldsymbol{b}_0$ ... momentum source - force

0 ... no momentum production

equilibrium equation $\quad \mathrm{D}_t(\rho_0\boldsymbol{v}) = \mathrm{Div}(\boldsymbol{P}) + \boldsymbol{b}_0$

compare

NEWTON'S THREE LAWS OF GRADUATION

First published in 1679, Isaac Newton's "***Procrastinare Unnaturalis Principia Mathematica***" is often considered one of the most important single works in the history of science. Its Second Law is the most powerful of the three, allowing mathematical calculation of the duration of a doctoral degree.

SECOND LAW

"The age, **a**, *of a doctoral process is directly proportional to the flexibility,* **f**, *given by the advisor and inversely proportional to the student's motivation,* **m**"

Mathematically, this postulate translates to:

$$age_{PhD} = \frac{flexibility}{motivation}$$

$$\mathbf{a} = \mathbf{F} / \mathbf{m}$$

$$\therefore \mathbf{F} = \mathbf{m}\,\mathbf{a}$$

This Law is a quantitative description of the effect of the forces experienced by a grad student. A highly motivated student may still remain in grad school given enough flexibility. As motivation goes to zero, the duration of the PhD goes to infinity.

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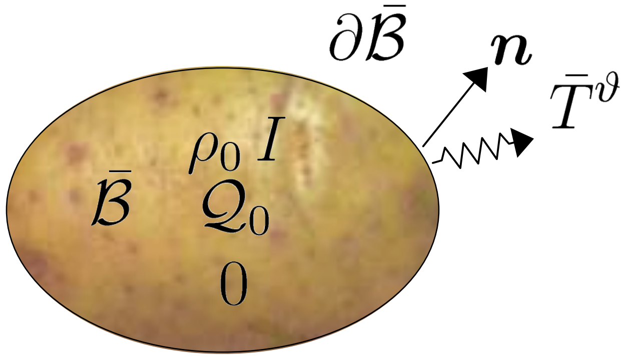
$$\mathrm{D}_t(\rho_0 \boldsymbol{v}) = \mathrm{Div}(\boldsymbol{P}) + \boldsymbol{b}_0$$

mass point

$$m\,\mathrm{D}_t \boldsymbol{v} = m\,\boldsymbol{a} = \boldsymbol{F}$$

balance of (internal) energy



$\rho_0\, I$ ... internal energy density

$\boldsymbol{Q}$ ... heat flux $\qquad -\boldsymbol{Q}\cdot\boldsymbol{n} = \bar{T}^{\vartheta}$

$\mathcal{Q}_0$ ... heat source

0 ... no heat production

energy equation $\mathrm{D}_t(\rho_0\, I) = \boldsymbol{P} : \mathrm{D}_t\boldsymbol{F} - \boldsymbol{v}\cdot\boldsymbol{b}_0 + \mathrm{Div}(-\boldsymbol{Q}) + \mathcal{Q}_0$

internal mechanical power $\qquad$ external thermal power