06 - kinematic equations -



continuum mechancis

continuum mechanics [kən'tın.ju.əm mə'kæn.iks] is a branch of physics (specifically mechanics) that deals with continuous matter. the fact that matter is made of atoms and that it commonly has some sort of heterogeneous microstructure is ignored in the simplifying approximation that physical quantities, such as energy and momentum, can be handled in the infinitesimal limit. differential equations can thus be employed in solving problems in continuum mechanics.



continuum mechanics

continuum mechanics [kən'tın.ju.əm mə'kæn.iks] is the branch of mechanics concerned with the stress in solids, liquids and gases and the deformation or flow of these materials. the adjective continuous refers to the simplifying concept underlying the analysis: we disregard the molecular structure of matter and picture it as being without gaps or empty spaces. we suppose that all the mathematical functions entering the theory are continuous functions. this hypothetical continuous material we call a continuum.

Malvern "Introduction to the mechanics of a continuous medium" [1969]

introduction

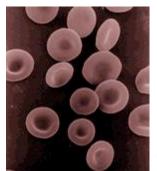
continuum mechanics

continuum hypothesis [kən'tın.ju.əm har'por0.ə.sıs] we assume that the characteristic length scale of the microstructure is much smaller than the characteristic length scale of the overall problem, such that the properties at each point can be understood as averages over a characteristic length scale

 $l^{
m micro} \ll l^{
m averg} \ll l^{
m conti}$

example: biomechanics

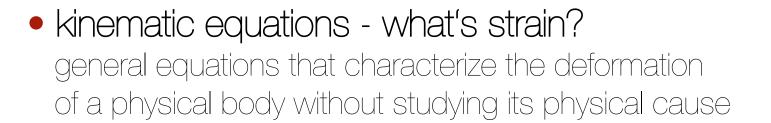
 $l^{\text{micro}} = l^{\text{cells}} \approx 10 \mu \text{m}$ $l^{\text{conti}} = l^{\text{tissue}} \approx 10 \text{cm}$



the continuum hypothesis can be applied when analyzing tissues

introduction





equations

- balance equations what's stress?
 general equations that characterize the cause of motion of any body
- constitutive equations how are they related? material specific equations that complement the set of governing equations

introduction

 $\epsilon = \frac{\Delta l}{l}$

 $\sigma = \frac{F}{A}$

 $\sigma = E \epsilon$

the potato 🥌 equations

- kinematic equations why not $\epsilon = \frac{\Delta l}{l}$? inhomogeneous deformation » non-constant finite deformation » non-linear inelastic deformation » growth tensor
 - $oldsymbol{F} =
 abla_X oldsymbol{arphi} \ oldsymbol{F} = oldsymbol{F}_{ ext{e}} \cdot oldsymbol{F}_{ ext{g}}$
- balance equations why not $\sigma = \frac{F}{A}$? $\text{Div}(\mathbf{P}) + \rho \mathbf{b}_0 = \mathbf{0}$ equilibrium in deformed configuration » multiple stress measures
- constitutive equations why not $\sigma = E \epsilon$? finite deformation » non-linear P = P(F)inelastic deformation » internal variables $P = P(\rho, F, F_g)$

introduction

kinematic equations

kinematic equations [kmə'mætik i'kwei.zəns] describe the motion of objects without the consideration of the masses or forces that bring about the motion. the basis of kinematics is the choice of coordinates. the 1st and 2nd time derivatives of the position coordinates give the velocities and accelerations. the difference in placement between the beginning and the final state of two points in a body expresses the numerical value of strain. strain expresses itself as a change in size and/or shape.



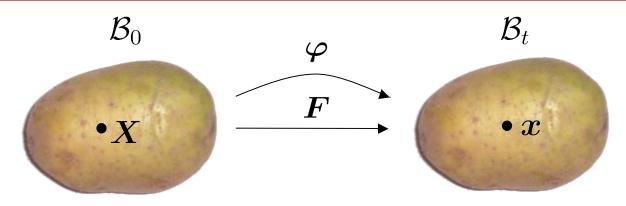
kinematic equations

kinematics [kmə'mætiks] is the study of motion per se, regardless of the forces causing it. the primitive concepts concerned are position, time and body, the latter abstracting into mathematical terms intuitive ideas about aggregations of matter capable of motion and deformation.

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Chadwick "Continuum mechanics" [1976]

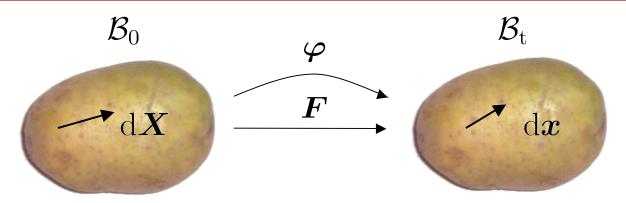
potato - kinematics



nonlinear deformation map φ $x = \varphi(X, t)$ with $\varphi : \mathcal{B}_0 \times \mathbb{R} \to \mathcal{B}_t$ spatial derivative of φ - deformation gradient

$$\mathrm{d}\boldsymbol{x} = \boldsymbol{F} \cdot \mathrm{d}\boldsymbol{X}$$
 with $\boldsymbol{F} : T\mathcal{B}_0 \to T\mathcal{B}_t$ $\boldsymbol{F} = \frac{\partial \varphi}{\partial \boldsymbol{X}}\Big|_{t \text{ fixed}}$

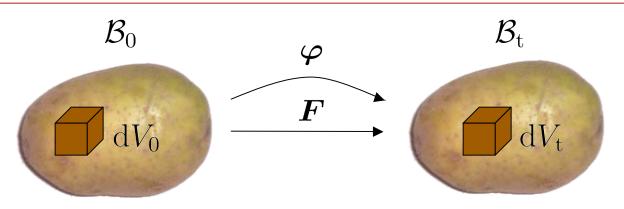
potato - kinematics



• transformation of line elements - deformation gradient F_{ij} $dx_i = F_{ij} dX_j$ with $F_{ij} : T\mathcal{B}_0 \to T\mathcal{B}_t$ $F_{ij} = \frac{\partial \varphi_i}{\partial X_j}\Big|_{t \text{ fixed}}$ • uniaxial tension (incompressible), simple shear, rotation

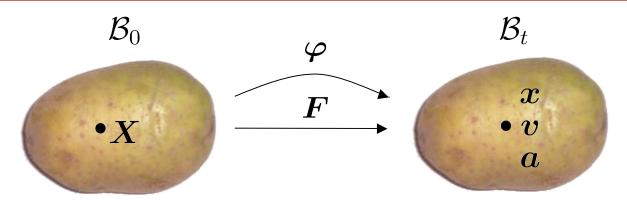
$$F_{ij}^{\text{uni}} = \begin{bmatrix} \alpha & 0 & 0 \\ 0 & \alpha^{-\frac{1}{2}} & 0 \\ 0 & 0 & \alpha^{-\frac{1}{2}} \end{bmatrix} F_{ij}^{\text{shr}} = \begin{bmatrix} 1 & \gamma & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix} F_{ij}^{\text{rot}} = \begin{bmatrix} \cos(\theta) & \sin(\theta) & 0 \\ -\sin(\theta) \cos(\theta) & 0 \\ 0 & 0 & 1 \end{bmatrix}$$

potato - kinematics of finite growth



• transformation of volume elements - determinant of F $dV_0 = dX_1 \cdot [dX_2 \times dX_3]$ $dV_t = dx_1 \cdot [dx_2 \times dx_3]$ $= det([dx_1, dx_2, dx_3])$ $= det([dX_1, dX_2, dX_3])$ $= det(F) det([dX_1, dX_2, dX_3])$ • changes in volume - determinant of deformation tensor J $dV_t = J dV_0$ J = det(F)

potato - kinematics



• temporal derivative of φ - velocity (material time derivative) $v = D_t \varphi = \frac{\partial \varphi}{\partial t}\Big|_{X \text{ fixed}}$ with $v : \mathcal{B}_0 \times \mathbb{R} \to \mathbb{R}^3$ • temporal derivative of v - acceleration $a = D_t v = \frac{\partial v}{\partial t}\Big|_{X \text{ fixed}} = \frac{\partial^2 \varphi}{\partial t^2}\Big|_{X \text{ fixed}}$ with $a : \mathcal{B}_0 \times \mathbb{R} \to \mathbb{R}^3$