

ME338A
CONTINUUM MECHANICS

lecture notes 09

tuesday, february 02nd, 2010

3.2 Balance of mass

total mass m of a body \bar{B}

$$m := \int_{\bar{B}} dm \quad (3.2.1)$$

mass density

$$\rho = \lim_{\Delta V \rightarrow 0} \frac{\Delta m}{\Delta V} = \frac{dm}{dV} \quad dm = \rho dV \quad (3.2.2)$$

$$m = \int_{\bar{B}} \rho dV \quad (3.2.3)$$

mass exchange of body \bar{B} with environment m^{sur} and m^{vol}

$$m^{\text{sur}} := \int_{\partial\bar{B}} r_n dA \quad m^{\text{vol}} := \int_{\bar{B}} \mathcal{R} dV \quad (3.2.4)$$

with contact mass flux $r_n = \mathbf{r} \cdot \mathbf{n}$ and mass source \mathcal{R}

3.2.1 Global form of balance of mass

"The time rate of change of the total mass m of a body \bar{B} is balanced with the mass exchange due to contact mass flux m^{sur} and the at-a-distance mass exchange m^{vol} ."

$$D_t m = m^{\text{sur}} + m^{\text{vol}} \quad (3.2.5)$$

and thus

$$D_t \int_{\bar{B}} \rho dV = \int_{\partial\bar{B}} r_n dA + \int_{\bar{B}} \mathcal{R} dV \quad (3.2.6)$$

3.2.2 Local form of balance of mass

modification of rate term $D_t m$

$$D_t m = D_t \int_{\bar{B}} \rho \, dV \stackrel{\bar{B}^{\text{fixed}}}{=} \int_{\bar{B}} D_t \rho \, dV \quad (3.2.7)$$

modification of surface term m^{sur}

$$m^{\text{sur}} = \int_{\partial \bar{B}} r_n \, dA \stackrel{\text{Cauchy}}{=} \int_{\partial \bar{B}} \mathbf{r} \cdot \mathbf{n} \, dA \stackrel{\text{Gauss}}{=} \int_{\bar{B}} \text{div}(\mathbf{r}) \, dV \quad (3.2.8)$$

and thus

$$\int_{\bar{B}} D_t \rho \, dV = \int_{\bar{B}} \text{div}(\mathbf{r}) \, dV + \int_{\bar{B}} \mathcal{R} \, dV \quad (3.2.9)$$

for arbitrary bodies $\bar{B} \rightarrow$ local form of mass balance

$$D_t \rho = \text{div}(\mathbf{r}) + \mathcal{R} \quad (3.2.10)$$

Classical continuum mechanics of closed systems

in classical closed system continuum mechanics (here), $\mathbf{r} = \mathbf{0}$ and $\mathcal{R} = 0$, such that the mass density ρ is constant in time

$$D_t \rho(\mathbf{x}, t) = 0 \quad \rho = \rho(\mathbf{x}) = \text{const} \quad (3.2.11)$$

typically $\mathbf{r} \neq \mathbf{0}$ and $\mathcal{R} \neq 0$ only in bio- or chemomechanics

3.3 Balance of linear momentum

total momentum \boldsymbol{p} of a body $\bar{\mathcal{B}}$

$$\boldsymbol{p} := \int_{\bar{\mathcal{B}}} \mathbf{D}_t \boldsymbol{u} \, d m = \int_{\bar{\mathcal{B}}} \boldsymbol{v} \, d m = \int_{\bar{\mathcal{B}}} \rho \boldsymbol{v} \, d V \quad (3.3.1)$$

momentum exchange of body $\bar{\mathcal{B}}$ with environment through contact forces $\boldsymbol{f}^{\text{sur}}$ and at-a-distance forces $\boldsymbol{f}^{\text{vol}}$

$$\boldsymbol{f}^{\text{sur}} := \int_{\partial \bar{\mathcal{B}}} \boldsymbol{t}_\sigma \, d A \quad \boldsymbol{f}^{\text{vol}} := \int_{\bar{\mathcal{B}}} \boldsymbol{b} \, d V \quad (3.3.2)$$

with contact/surface force $\boldsymbol{t}_\sigma = \boldsymbol{\sigma}^t \cdot \boldsymbol{n}$ and volume force \boldsymbol{b}

3.3.1 Global form of balance of momentum

"The time rate of change of the total momentum \boldsymbol{p} of a body $\bar{\mathcal{B}}$ is balanced with the momentum exchange due to contact momentum flux / surface force $\boldsymbol{f}^{\text{sur}}$ and the at-a-distance momentum exchange / volume force $\boldsymbol{f}^{\text{vol}}$."

$$\mathbf{D}_t \boldsymbol{p} = \boldsymbol{f}^{\text{sur}} + \boldsymbol{f}^{\text{vol}} \quad (3.3.3)$$

and thus

$$\mathbf{D}_t \int_{\bar{\mathcal{B}}} \rho \boldsymbol{v} \, d V = \int_{\partial \bar{\mathcal{B}}} \boldsymbol{t}_\sigma \, d A + \int_{\bar{\mathcal{B}}} \boldsymbol{b} \, d V \quad (3.3.4)$$

3.3.2 Local form of balance of momentum

modification of rate term $\mathbf{D}_t \boldsymbol{p}$

$$\mathbf{D}_t \boldsymbol{p} = \mathbf{D}_t \int_{\bar{\mathcal{B}}} \rho \boldsymbol{v} \, d V \stackrel{\bar{\mathcal{B}}^{\text{fixed}}}{=} \int_{\bar{\mathcal{B}}} \mathbf{D}_t (\rho \boldsymbol{v}) \, d V \quad (3.3.5)$$

modification of surface term f^{sur}

$$f^{\text{sur}} = \int_{\partial\bar{\mathcal{B}}} \mathbf{t}_\sigma \, dA \stackrel{\text{Cauchy}}{=} \int_{\partial\bar{\mathcal{B}}} \boldsymbol{\sigma}_t \cdot \mathbf{n} \, dA \stackrel{\text{Gauss}}{=} \int_{\bar{\mathcal{B}}} \text{div}(\boldsymbol{\sigma}^t) \, dV \quad (3.3.6)$$

and thus

$$\int_{\bar{\mathcal{B}}} \mathbf{D}_t(\rho \mathbf{v}) \, dV = \int_{\bar{\mathcal{B}}} \text{div}(\boldsymbol{\sigma}^t) \, dV + \int_{\bar{\mathcal{B}}} \mathbf{b} \, dV \quad (3.3.7)$$

for arbitrary bodies $\bar{\mathcal{B}} \rightarrow$ local form of momentum balance

$$\mathbf{D}_t(\rho \mathbf{v}) = \text{div}(\boldsymbol{\sigma}^t) + \mathbf{b} \quad (3.3.8)$$

3.3.3 Reduction with lower order balance equations

balance equations are typically modified with the help of lower order balance equations

$$\mathbf{D}_t(\rho \mathbf{v}) = \mathbf{v} \mathbf{D}_t \rho + \rho \mathbf{D}_t \mathbf{v} = \text{div}(\boldsymbol{\sigma}^t) + \mathbf{b} \quad (3.3.9)$$

with balance of mass multiplied by velocity \mathbf{v}

$$\mathbf{v} \mathbf{D}_t \rho = \mathbf{v} \text{div}(\mathbf{r}) + \mathbf{v} \mathcal{R}_0 = \text{div}(\mathbf{v} \otimes \mathbf{r}) - \nabla \mathbf{v} \cdot \mathbf{r} + \mathbf{v} \mathcal{R}_0 \quad (3.3.10)$$

local momentum balance in reduced format

$$\rho \mathbf{D}_t \mathbf{v} = \text{div}(\boldsymbol{\sigma}^t - \nabla \mathbf{v} \otimes \mathbf{r}) + \mathbf{b} + \nabla \mathbf{v} \cdot \mathbf{r} - \mathbf{v} \mathcal{R}_0 \quad (3.3.11)$$